

Probing TeV physics using neutron decays and nEDM

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PNDME Collaboration: PRD 85 (2012) 054512 ; PRD 89 (2014) 094502

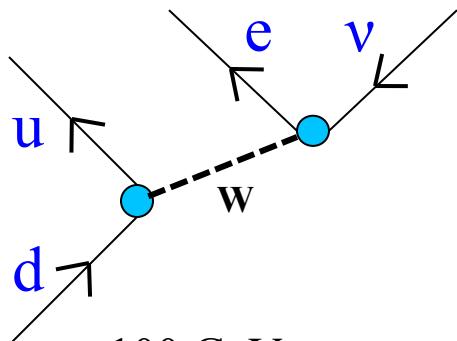
- Tanmoy Bhattacharya, Rajan Gupta, Boram Yoon (LANL)
- Saul Cohen, Huey-Wen Lin (UW)
- Anosh Joseph (DESY, Zeuthen)

nEDM (LANL-RBC-UKQCD) Collaboration:

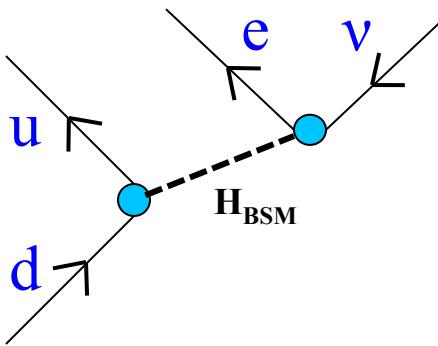
- Tanmoy Bhattacharya, Vincenzo Cirigliano, Rajan Gupta, Boram Yoon (LANL)
- Michael Abramczyk and Tom Blum (RBC, U of Connecticut)
- Taku Izubuchi, Amarjit Soni, Shigemi Ohta, Christoph Lehner and Sergey Syritsyn (RBC, BNL)

Probing New Interactions: $M_{\text{BSM}} \gg M_W \gg 1 \text{ GeV}$

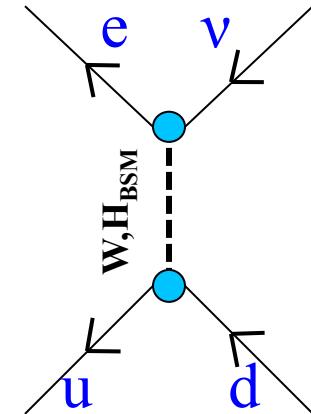
Many BSM possibilities for Scalar & Tensor: Higgs-like, leptoquark, loop effects, ...



100 GeV
V-A Weak decay
 e_L with a ν_L

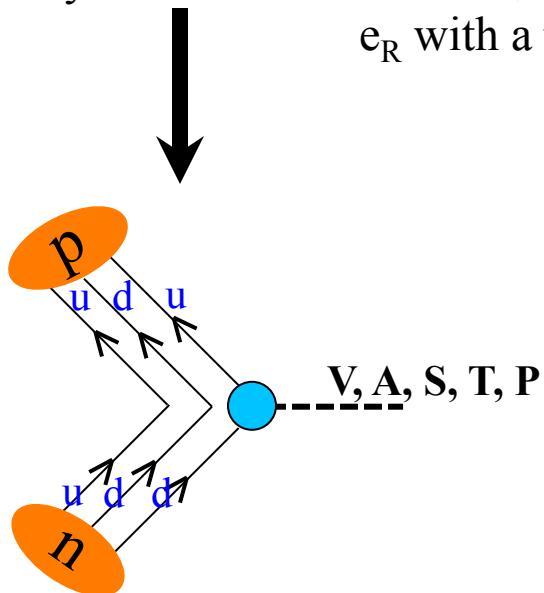


TeV
Novel S,T Interactions
 e_R with a ν_L



LHC

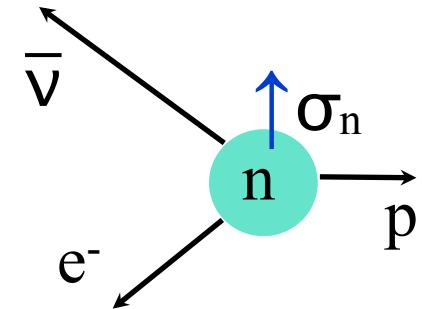
Neutron
Decay



Effective Theory @ ~ 2 GeV
V-A (g_A, g_V) Weak interactions
S, T (g_S, g_T) New Interactions

[Ultra] Cold Neutron Decay: Parameters sensitive to new physics

Neutron decay can be parameterized as



$$d\Gamma \propto F(E_e) \left[1 + b \frac{m_e}{E_e} + \left(B_0 + B_1 \frac{m_e}{E_e} \right) \frac{\vec{\sigma}_n \cdot \vec{p}_\nu}{E_\nu} + \dots \right]$$

b: Deviations from the leading order electron spectrum:
Fierz interference term

B₁: Energy dependent part of antineutrino correlation
with neutron spin

Relating b and B_1 to ME & BSM couplings

$$H_{eff} \supset G_F \left[\varepsilon_S \boxed{\bar{u}d} \times \bar{e}(1 - \gamma_5)v_e + \varepsilon_T \boxed{\bar{u}\sigma_{\mu\nu}d} \times \bar{e}\sigma^{\mu\nu}(1 - \gamma_5)v_e \right]$$

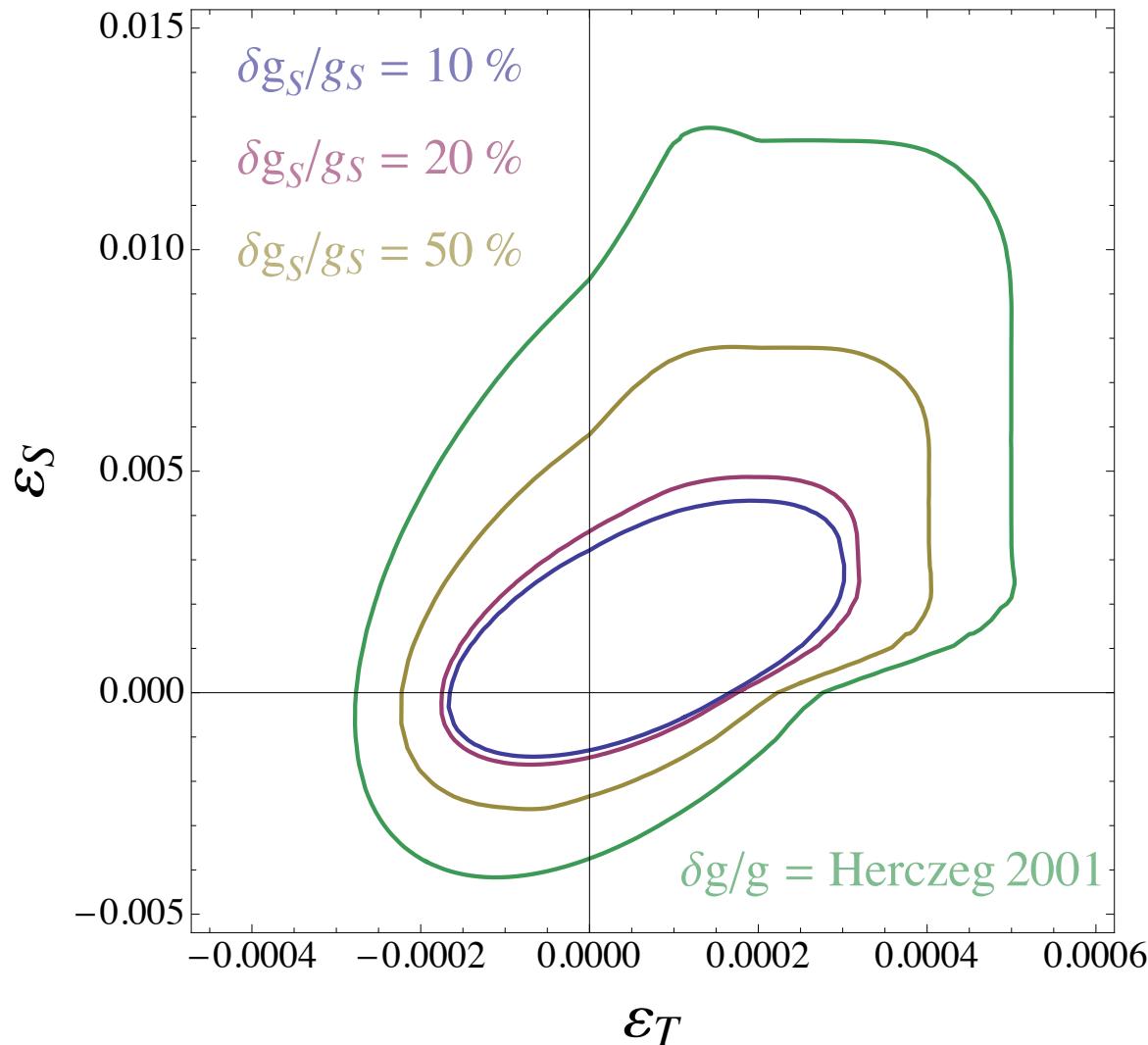
$$g_S \sim \langle p | \bar{u}d | n \rangle \quad g_T \sim \langle p | \bar{u}\sigma_{\mu\nu}d | n \rangle$$

Linear order relations from $n \rightarrow p e \bar{\nu}$ decay

$$b^{BSM} \approx 0.34g_S\varepsilon_S - 5.22g_T\varepsilon_T$$

$$b_v^{BSM} \equiv B_1^{BSM} = E_e \frac{\partial B^{BSM}(E_e)}{\partial m_e} \approx 0.44g_S\varepsilon_S - 4.85g_T\varepsilon_T$$

Target Precision for g_S , g_T : 10-20%



Expt. input

$|B_1 - b| < 10^{-3}$
 $|b| < 10^{-3}$
 $b_{0+} = 2.2 (4.3) * 10^{-3}$

Allowed region in $[\epsilon_S, \epsilon_T]$ (90% contours)

Clover on 2+1+1 flavor HISQ lattices: ~1000 configs

- m_s set to its physical value using $M_{\bar{S}S}$

$a(\text{fm})$	m_l/m_s	Lattice Volume	$M_\pi L$	$M_\pi (\text{MeV})$	Configs. X sources
0.12	0.2	$24^3 \times 64$	4.54	305	1013×8
0.12	0.1	$24^3 \times 64$	3.22	217	1000×12
0.12	0.1	$32^3 \times 64$	4.3	217	958×8
0.12	0.1	$40^3 \times 64$	5.36	217	1010×8
0.09	0.2	$32^3 \times 96$	4.5	313	881×8
0.09	0.1	$48^3 \times 96$	4.71	220	890×8
0.09	0.035	$64^3 \times 96$	3.66	130	883×4
0.06	0.2	$48^3 \times 144$	4.51	320	865×4
0.06	0.1	$64^3 \times 144$	4.25	229	200×4

Observations and Lessons Learned

- Exceptional configurations: Clover-on-HISQ with HYP smearing:
 - Exist on the $a=0.12\text{fm}$ lattices with $M_\pi=130\text{MeV}$
 - None found for $a=0.09$ and 0.06 lattices
- Statistics: $\sigma(g_S) \sim 5 \sigma(g_A)$ [or $5 \sigma(g_T)$]
Need $O(15000)$ independent measurement (Configs \times Sources)
- Excited state contamination is significant
Need multiple T_{sep} with $T_{\text{sep}} > 1.2\text{fm}$
 $T_{\text{sep}} > 1.2\text{fm}$ and fits including at least one excited state
- Renormalization (RI-sMOM): Smearing introduces artifacts.
Need a prescription with a well-defined continuum limit

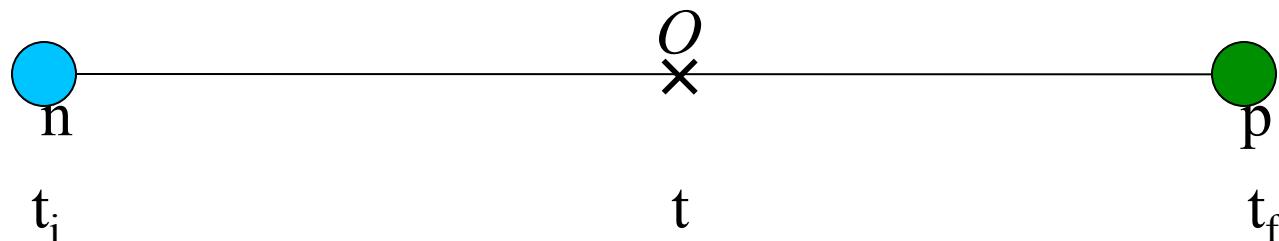
Reducing excited state contamination

Simultaneous fit to all $\Delta t = t_{sep} = t_f - t_i$

Assuming 1 excited state, the 3-point function is given by

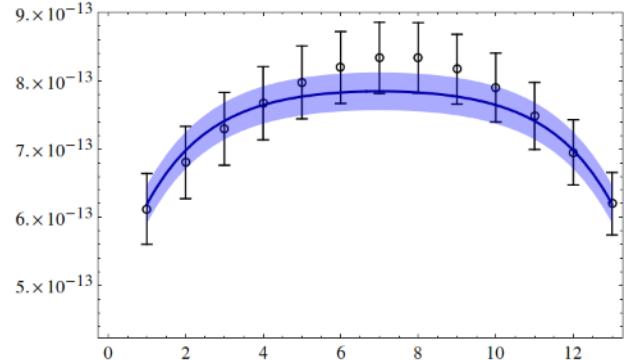
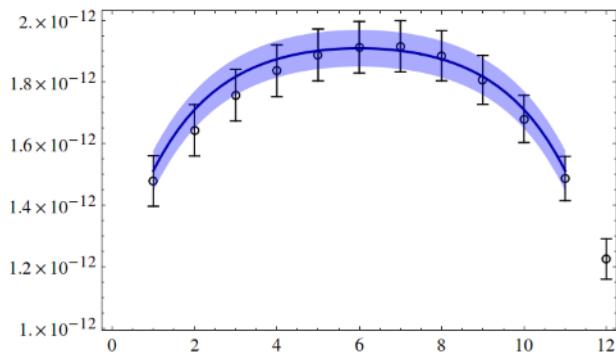
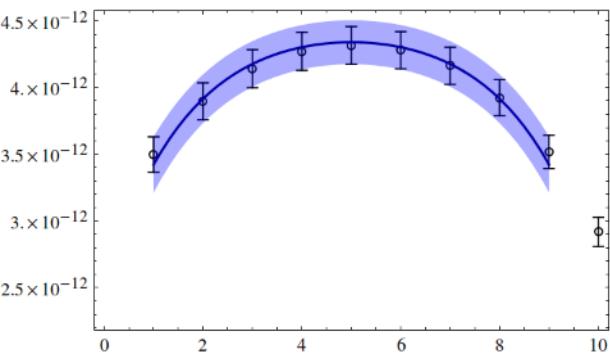
$$\Gamma^3(t_f, t, t_i) = |A_0|^2 \langle 0|O|0\rangle e^{-M_0 \Delta t} + |A_1|^2 \langle 1|O|1\rangle e^{-M_1 \Delta t} + A_0 A_1^* \langle 0|O|1\rangle e^{-M_0 \Delta t} e^{-M_1 (\Delta t - t)} + A_0^* A_1 \langle 1|O|0\rangle e^{-M_1 \Delta t} e^{-M_0 (\Delta t - t)}$$

Where M_0 and M_1 are the masses of the ground & excited state and A_0 and A_1 are the corresponding amplitudes.

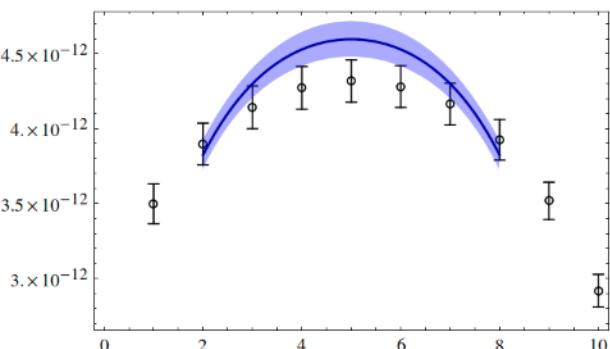


Simultaneous fit to all t_{sep}

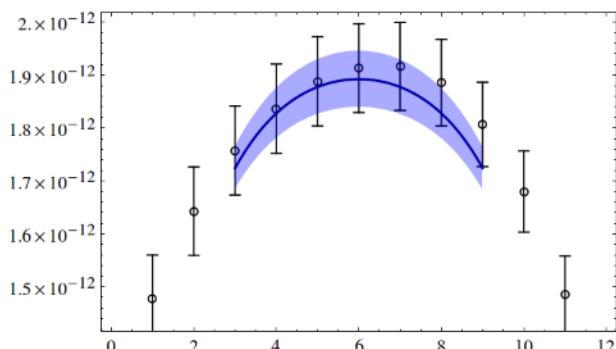
Data for g_S on the $M_\pi=220$ MeV ensemble at $a=0.09\text{fm}$



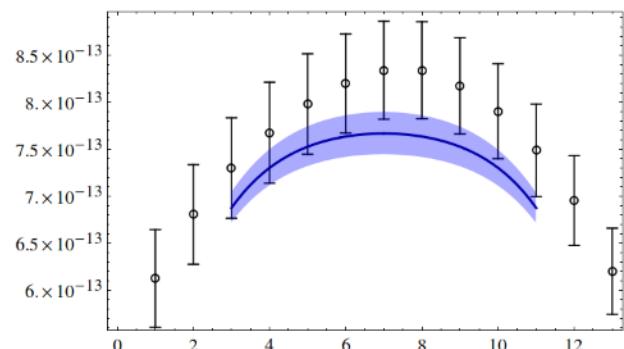
Excluding the $\langle 1|O|1 \rangle$ Term



$t_{sep}=10$



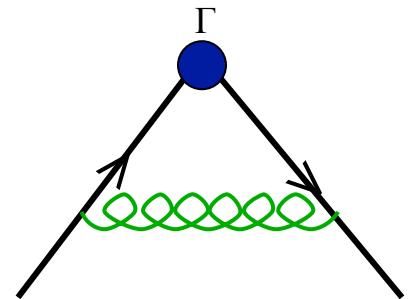
$t_{sep}=12$



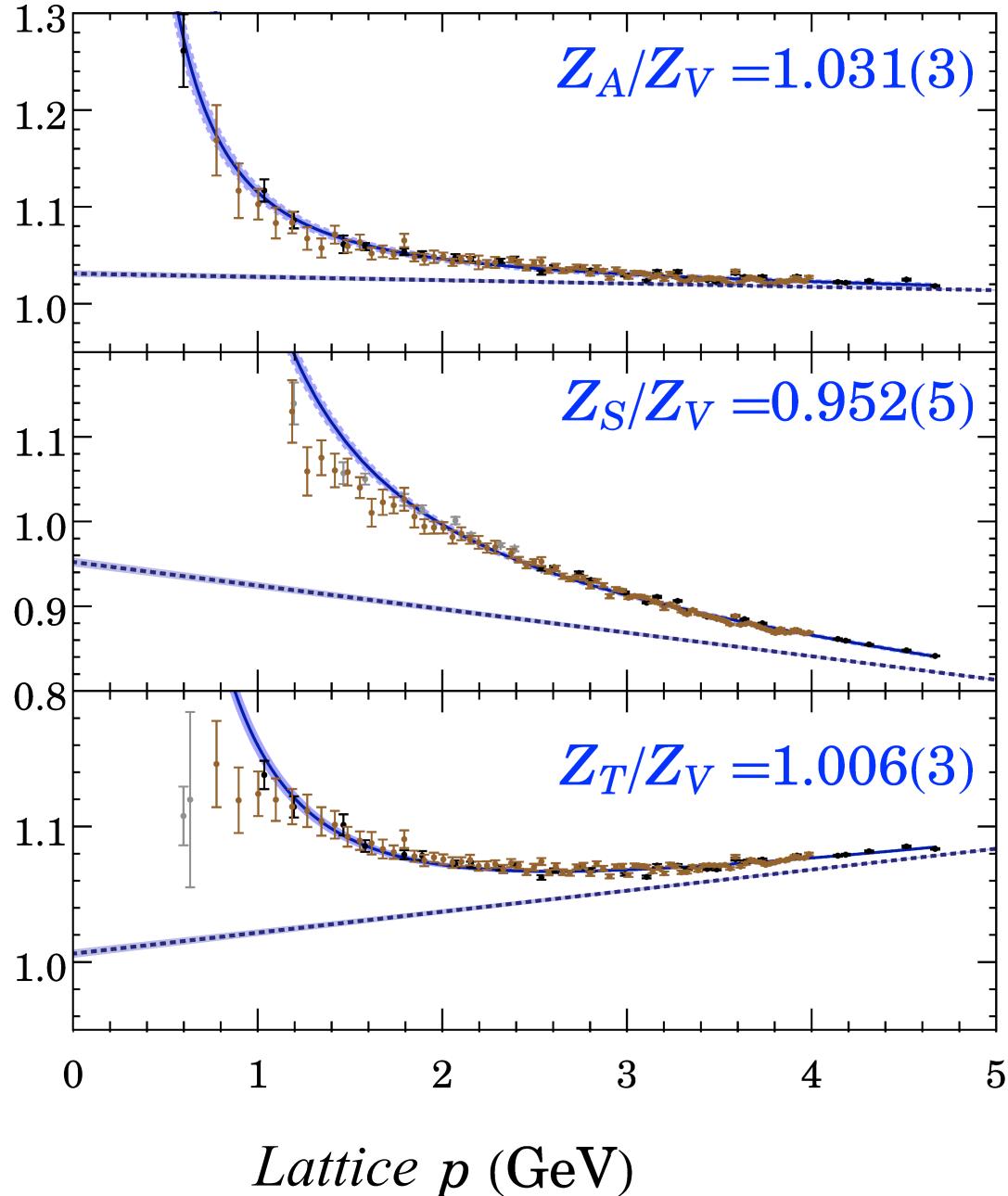
$t_{sep}=14$

Renormalization of bilinear operators

- Non-perturbative renormalization Z_Γ using the RI-sMOM scheme
 - Need quark propagator in momentum space
- Assumption: there exists a window
 - $\Lambda_{QCD} \ll p \ll \pi/a$
- HYP Smearing introduces artifacts
 - Gluon momentum above $\sim 1/a$ are an average
 - A window may no longer exist on coarse lattices
- No detectable dependence of Z 's on m_q



$$\frac{Z_{A,S,T}}{Z_V} \left(\overline{MS}, 2 \text{ GeV} \right)$$



Fit data to: $A/p + Z + Cp$
in the range $\{1 < p < 4 \text{ GeV}\}$
Or
Choose Z at $p^2 = 5 \text{ GeV}^2$
& errors from $\{4 < p^2 < 6 \text{ GeV}^2\}$

Renormalized Charges

$$\frac{Z_{A,S,T}}{Z_V} \times \frac{g_{A,S,T}}{g_V}$$

Use Ward identity $Z_V g_V = 1$

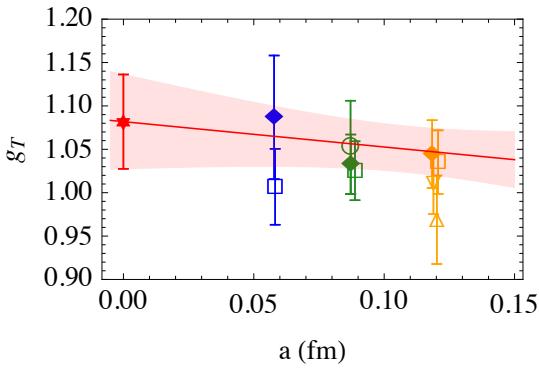
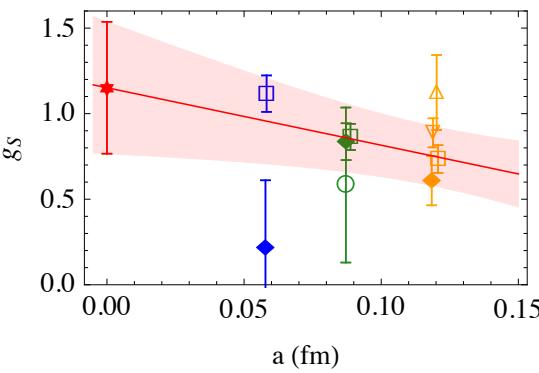
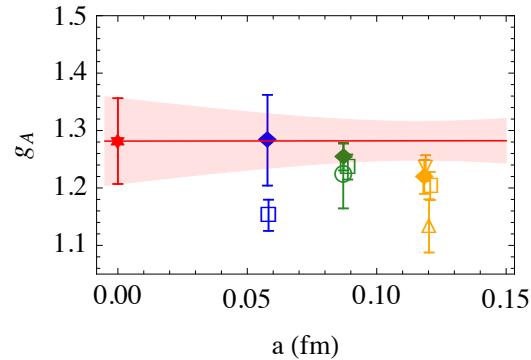
Extrapolations in a, M_π^2, L

$$g(a, M_\pi, L) = g + A a + B M_\pi^2 + C e^{-M_\pi L} + \dots$$

We use the lowest order corrections when fitting 9 points

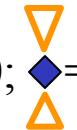
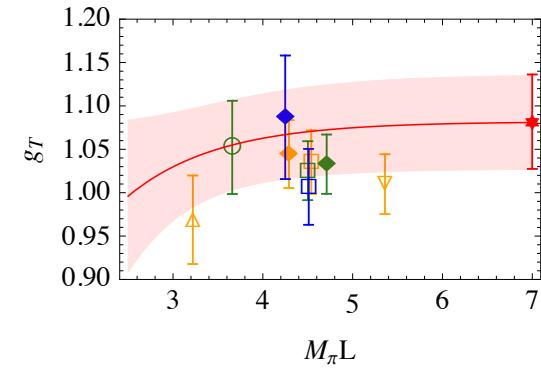
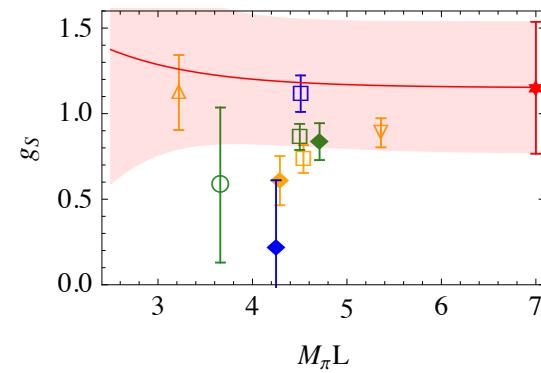
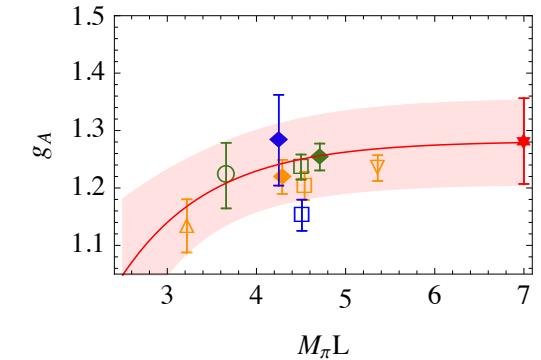
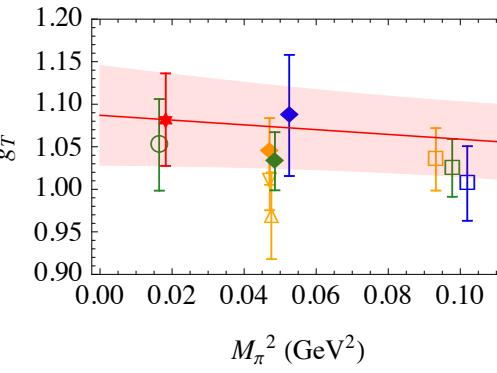
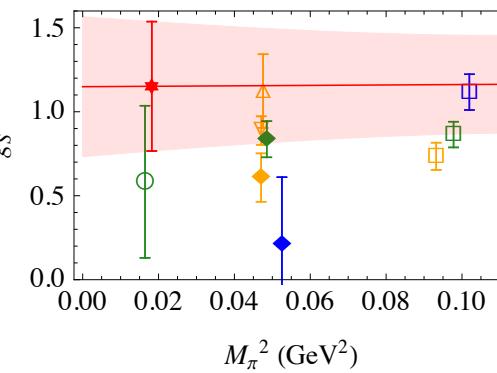
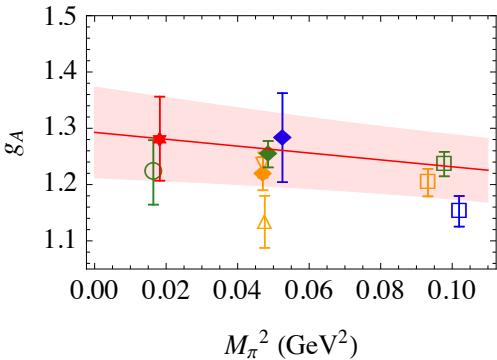
- Lattice spacing a
- Dependence on quark mass $m_q \sim M_\pi^2$
- Finite volume

★: Preliminary Results at the physical point



Orange: $a=0.12$; Green: $a=0.09$; Blue: $a=0.06$ fm

$\square=310$; $\diamond=220$; $\circ=130$ MeV



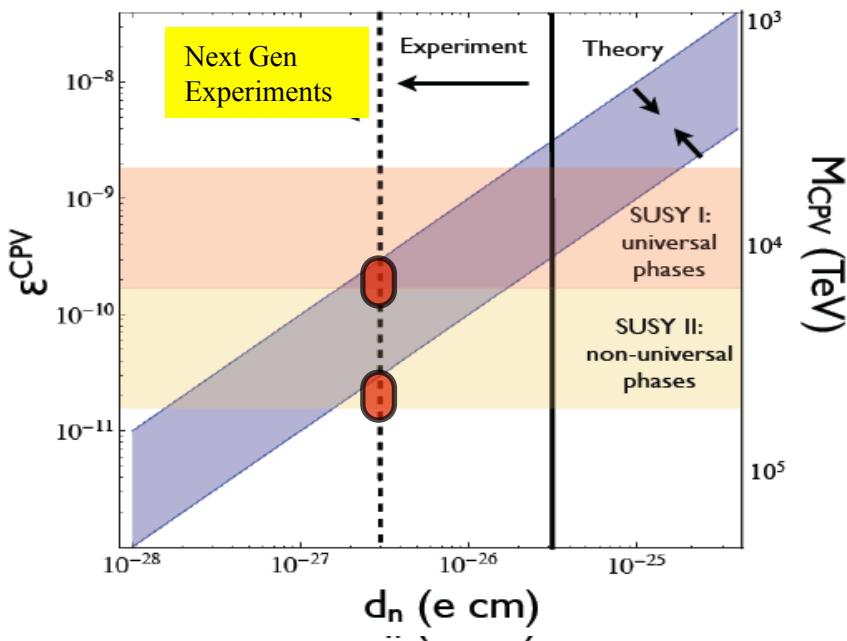
Towards Physical Estimates

- g_T : $1.08(5)$ [preliminary]
 - Estimate of Z_T is reliable
 - Small dependence on $a, M_\pi^2, M_\pi L$
- g_A :
 - Estimate of Z_A is reliable
 - $\exp\{-M_\pi L\}$ is the largest effect → Expt. result
- g_S :
 - Statistical errors are large
 - Z_S is not well-determined
 - Extrapolations in $a, M_\pi^2, M_\pi L$ are not stable

CP Violation, Baryogenesis, nEDM

CP Violation, Baryogenesis, nEDM

- CPV in SM is too small to explain Baryogenesis
- BSM theories have new sources of larger CPV that can explain baryogenesis
- These novel CPV may also give rise to larger nEDM ($d_n \sim 10^{-27}$ e-cm)
- Next generation nEDM will push limit from ($d_n \sim 3 \times 10^{-26}$ e-cm) to ($d_n \sim 10^{-28}$ e-cm)
- Low energy effective theory at ~ 2 GeV: Two leading dimension-5 CPV operators
- To use experimental value of d_n to constrain BSM, need high precision ME

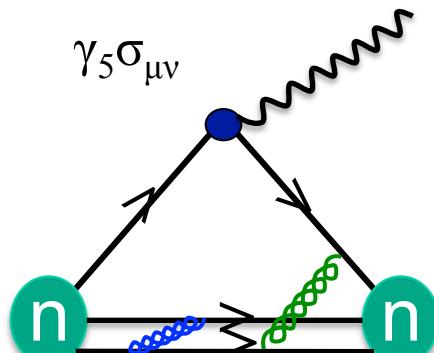


EXAMPLE: 10^{-27} e-cm sensitivity in nEDM experiments requires $O(1)$ accuracy in ME to constrain BSM

Current uncertainty in ME is $O(10)$
[Engels, Ramsey-Musolf, Kolck, Prog. Part. Nucl. Phys., 71]

ME of novel CP violating operators: nEDM

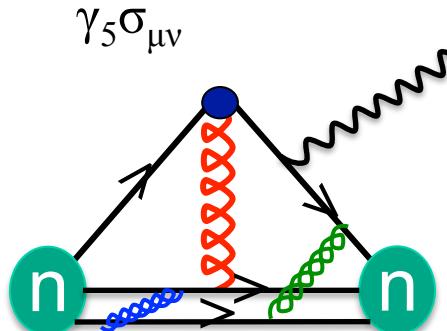
Quark-EDM



γ attaches to the vertex

$$\bar{q} \sigma_{\mu\nu} \gamma_5 q F^{\mu\nu}$$

Chromo-EDM



- 4-pt function as γ can attach to any quark line
- Gluon free end can attach to any quark line

$$\bar{q} \sigma_{\mu\nu} \gamma_5 q \lambda^a G_a^{\mu\nu}$$

- Formulation of the problem
- Operator mixing and renormalization
- Signal in disconnected diagrams
- Formulating lattice calculation of chromo EDM, a 4-point function

Mixing between D-5 CP violating operators in MS schemes

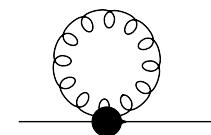
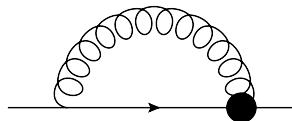
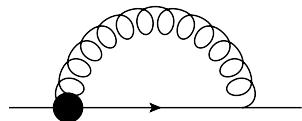
	E	C	P_{EE}	$\partial \cdot A_E$	A_G	$\partial^2 P$	mP_E	$m\partial \cdot A$	$m^2 P$	mG
E	.	0	0	0	0	0	0	0	0	0
C
P_{EE}	0	0	.	.	.	0	.	0	0	0
$\partial \cdot A_E$	0	0	0	.	0	.	0	.	0	0
A_G	0	0	.	.	.	0	.	0	0	0
$\partial^2 P$	0	0	0	0	0	.	0	0	0	0
mP_E	0	0	0	0	0	0	.	0	0	0
$m\partial \cdot A$	0	0	0	0	0	0	0	.	0	0
$m^2 P$	0	0	0	0	0	0	0	0	.	0
mG	0	0	0	0	0	0	0	.	0	.

The operators are defined as

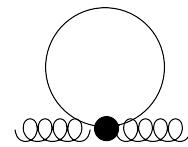
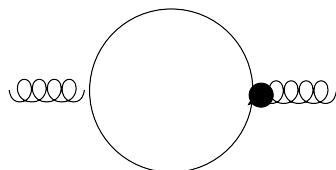
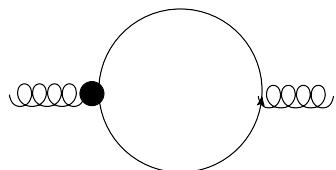
$$\begin{aligned}
 E \quad e &\equiv \bar{\psi} \sigma \cdot \tilde{F} \psi, \\
 C \quad g &\equiv \bar{\psi} \sigma \cdot \tilde{G} \psi, \\
 P_{EE} &\equiv \bar{\psi} (i \not{D} - m) \gamma_5 (i \not{D} - m) \psi, \\
 A_E^\mu &\equiv \bar{\psi} \gamma_\mu \psi + h.c., \\
 A_G &\equiv i \bar{\psi} (i \not{D} + m) \not{\partial} \psi + h.c., \\
 P &\equiv \bar{\psi} \gamma_5 \psi, \\
 P_E &\equiv \bar{\psi} (i \not{D} + m) \gamma_5 \psi + h.c., \\
 A^\mu &\equiv \partial_\mu \bar{\psi} \gamma^\mu \gamma_5 \psi, \\
 G &= G_{\mu\nu} \tilde{G}^{\mu\nu}
 \end{aligned}$$

- Entries marked 0 are prohibited by the structure of the operators.
- Only E , C , P , and G contribute to physical ('on-shell') matrix elements at zero momentum
- Rest need to be evaluated to connect the minimal subtraction scheme to a regularization independent off-shell scheme like RI-sMOM.

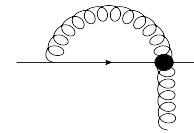
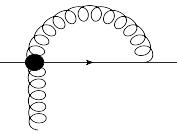
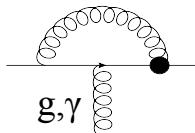
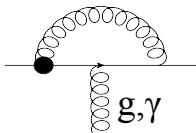
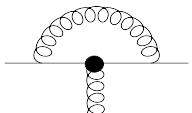
Feynman diagrams for CEDM



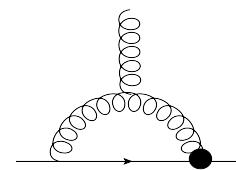
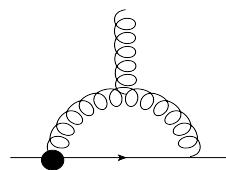
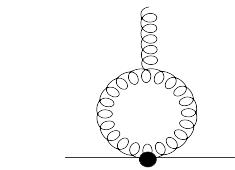
Diagrams contributing to the quark 2-pt function



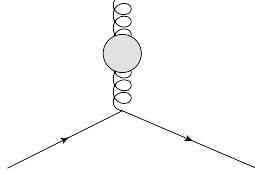
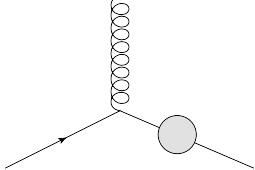
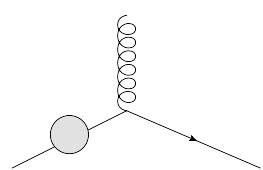
Diagrams contributing to the gluon 2-pt function



IPI diagrams contributing to the quark 3-pt function



Non IPI diagrams contributing to the quark 3-pt function



Summary

- Calculation of g_A , g_S , g_T on track
 - Need higher statistics to get g_S with 10%
 - Need further finite volume study especially for g_A
 - $g_T = 1.08(5)$ [preliminary]
- QEDM calculations are on track
 - DWF on DWF [See talk by Michael Abramczyk]
 - Disconnected diagrams [See talk by Boram Yoon]
- CEDM calculation is hard
 - Mixing and renormalization: perturbative calculations are nearing completion
 - We are formulating the lattice calculation of the ME

Acknowledgements

- Computing resources from
 - USQCD
 - XSEDE
 - LANL
- 2+1+1 HISQ lattices generated by the MILC collaboration
- Computer code uses CHROMA library
- Supported by DOE and LANL-LDRD